

Planetenbewegung in Sternsystemen

The Effect of Resonances

Part 1

Topics overview

1. Definition and examples of “resonances”
2. Disturbing function
3. Mean-motion resonance (MMR)
4. Secular resonance (SR)
5. Kozai-Lidov resonance
6. Evection resonance
7. Other resonances

Recap

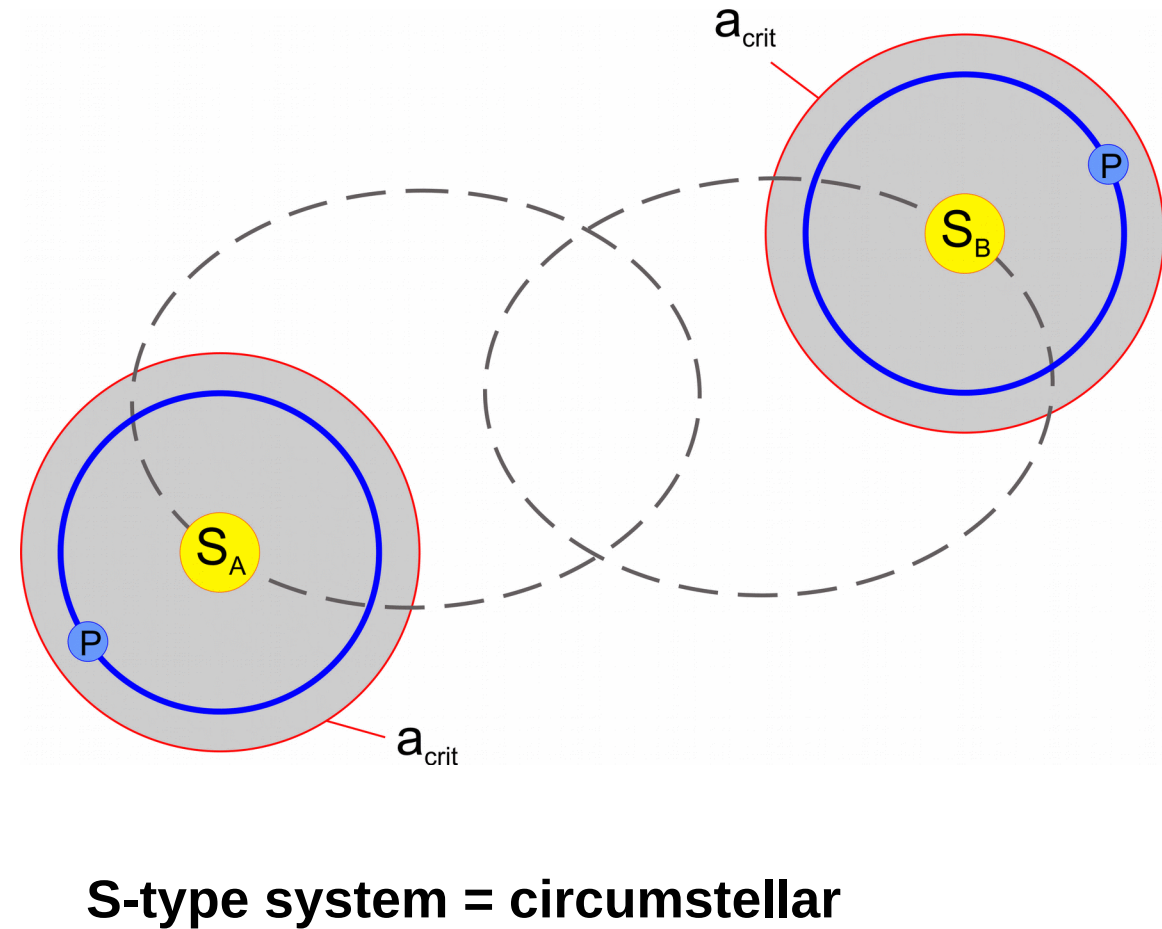
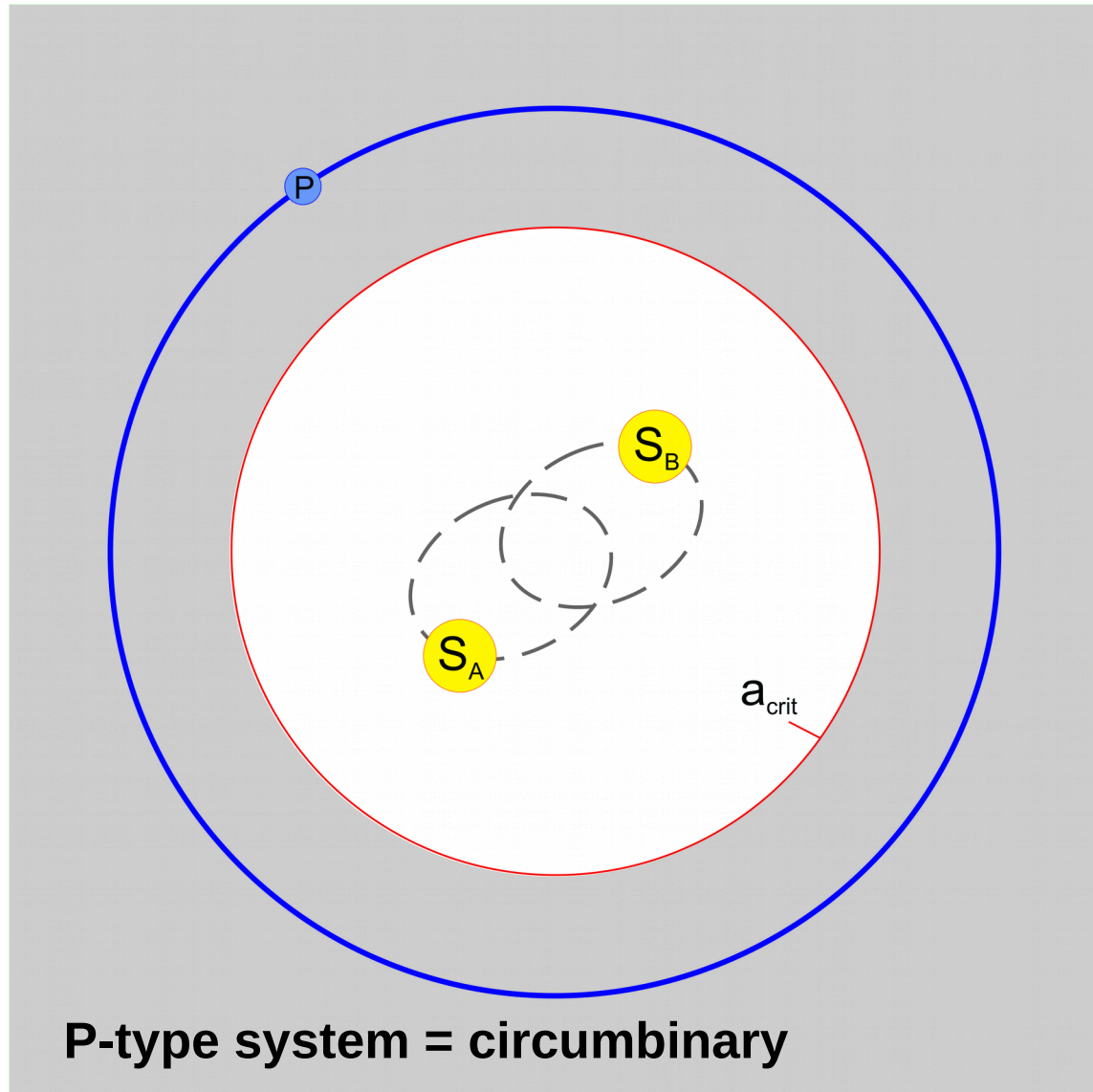
- Habitability → long-term stability → resonances
- Multi-stellar systems with exoplanets: 2, 3, 4 stars

https://en.wikipedia.org/wiki/Star_system

<http://www.univie.ac.at/adg/schwarz/multiple.html>

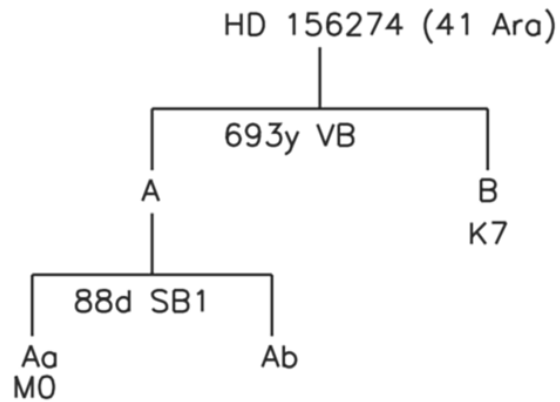
- Binary stars and hierarchical systems

Binary stars – domains of regular motion

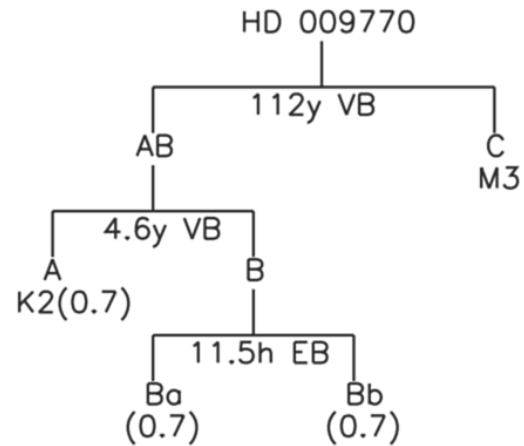


Higher multiplicity stellar systems

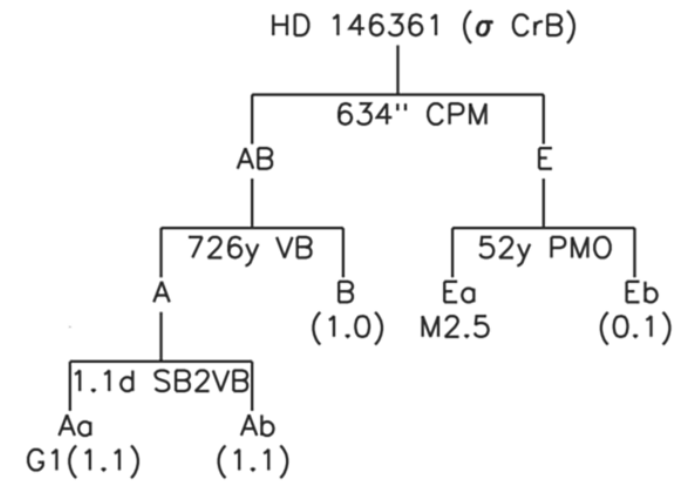
Triple system



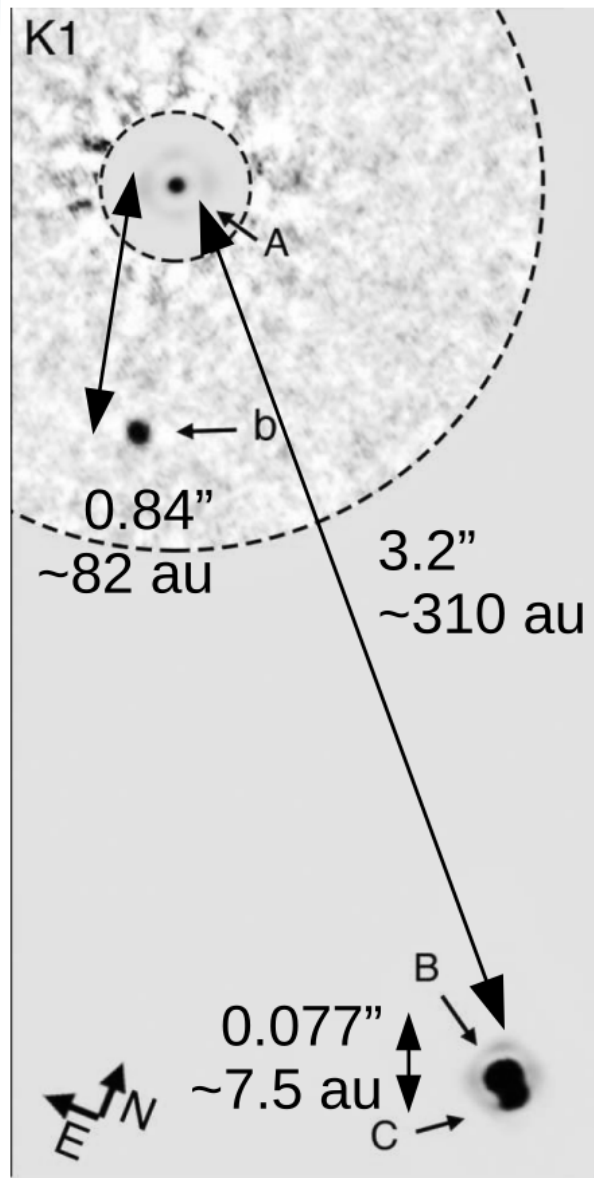
Quadruple system



Quintuple system



Higher multiplicity – HD 131399



- Direct imaging (NIR adaptive optics + coronagraph)
- Dynamical configuration: Ab + BC
- Star A: A1V ($1.82 M_{\odot}$)
- Star B: G ($0.96 M_{\odot}$)
- Star C: K ($0.6 M_{\odot}$)

Wagner,+ (2016)

1. Resonances

Resonances

Definition

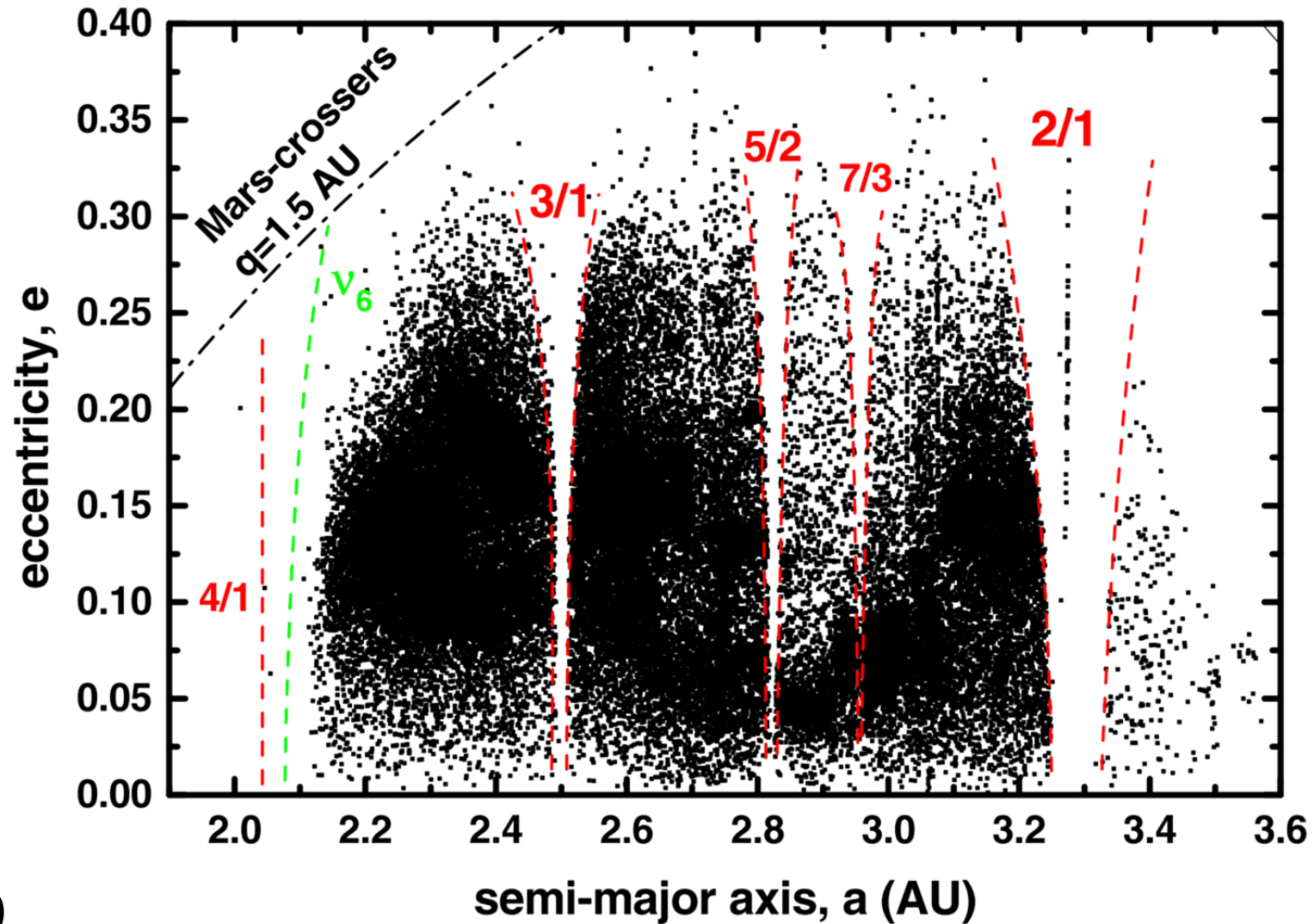
Two frequencies f_1 and f_2 are in resonance if their ratio can be expressed as a rational number

$$\left| \frac{f_1}{f_2} \right| = \frac{p}{q}, \quad p, q \in \mathbb{N}$$

Types of resonances

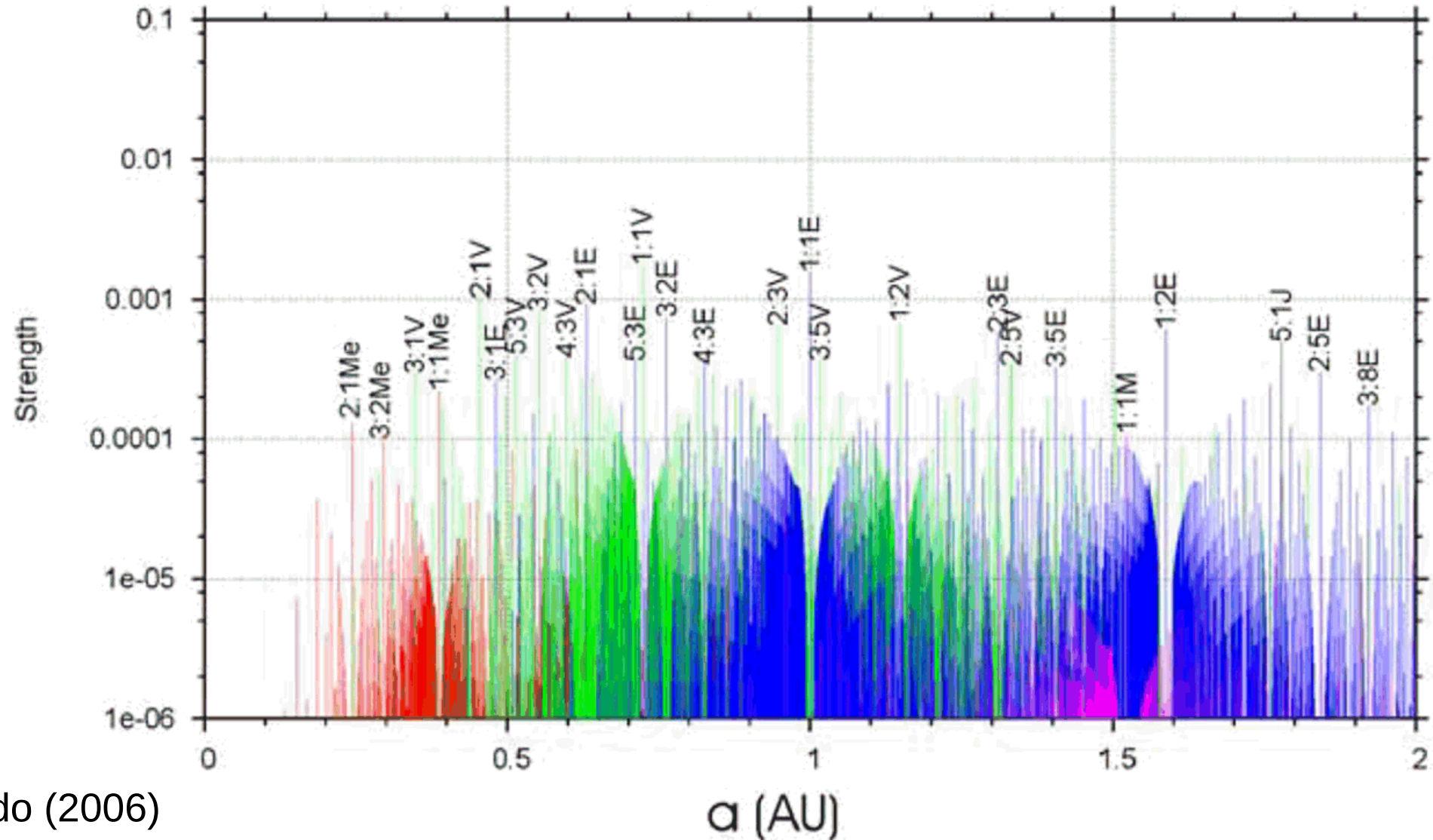
- **Mean-motion resonance (MMR)**
 - Orbital frequencies of two bodies
- **Three-body resonance**
 - Orbital frequencies of three bodies
- **Secular resonance (SR)**
 - Precession frequencies (perihelion, node) of ≥ 2 bodies
- **Spin-orbit resonance**
 - Orbital and spin frequency of same body
- **Gravitational resonance**
 - Orbital and spin frequency of different bodies

Examples – MMR



Tsiganis (2010)

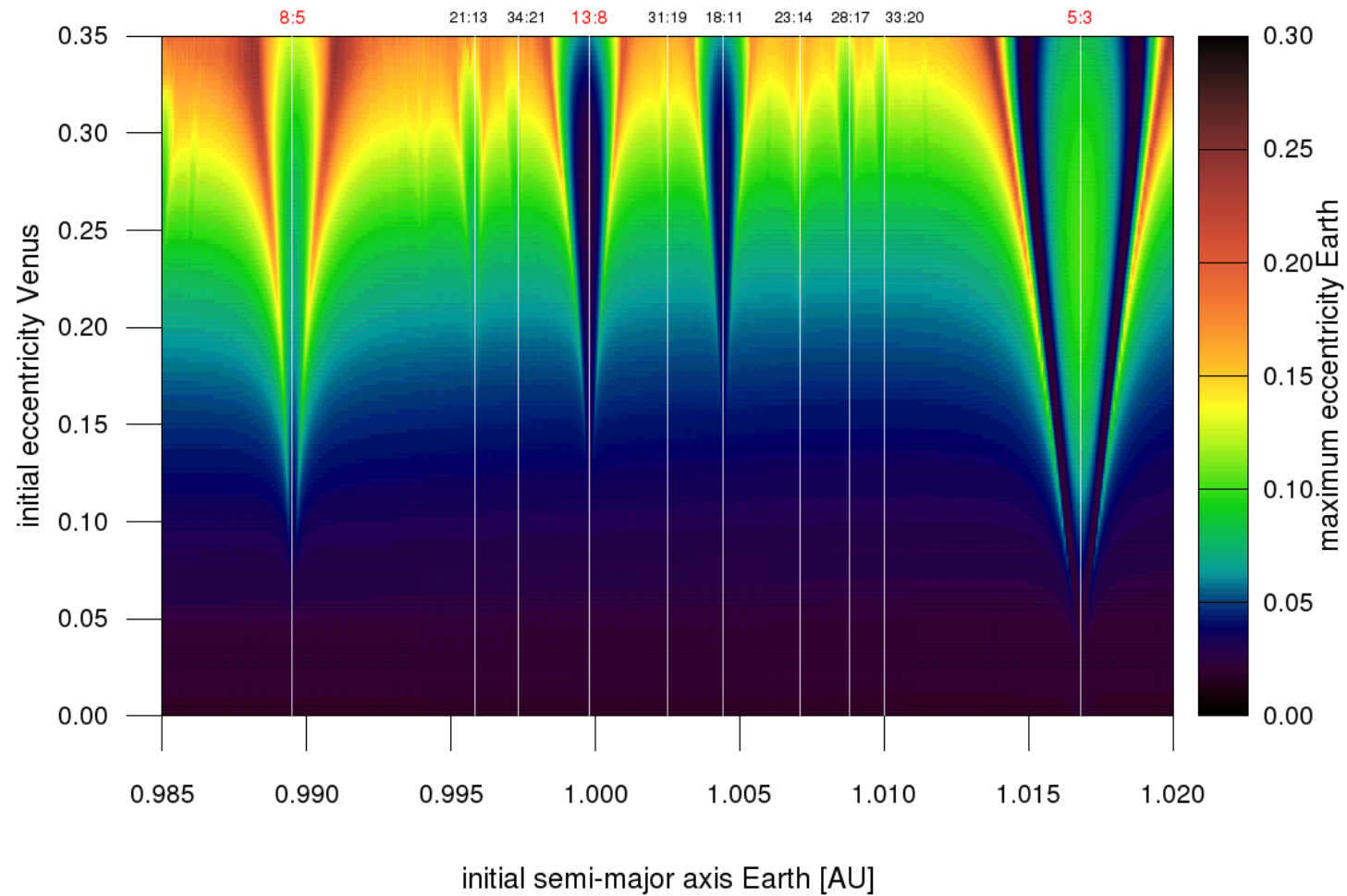
Examples - MMR



Gallardo (2006)

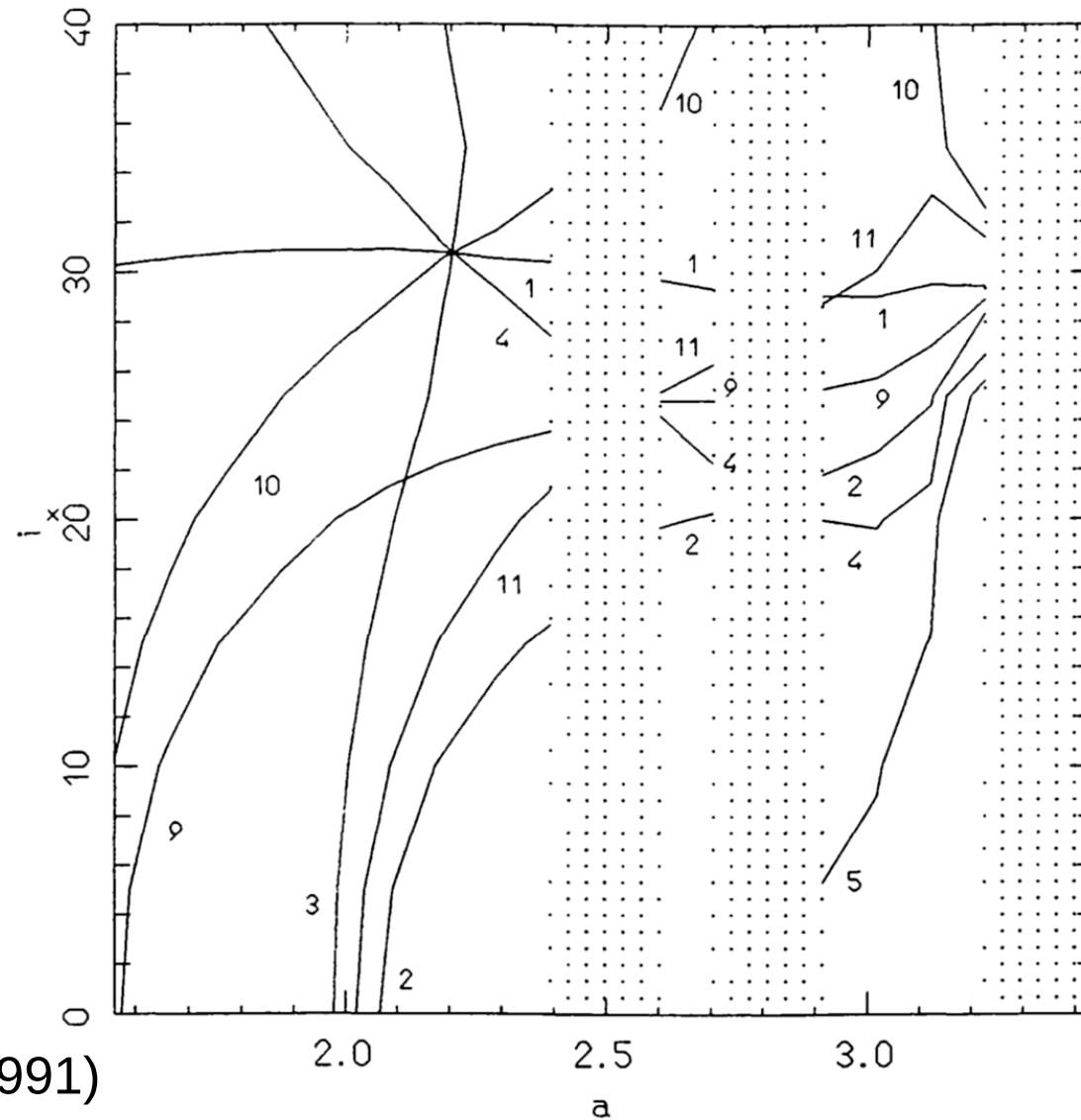
Examples – MMR

Overview of some MMR between Venus-Earth



Bazsó,+ (2010)

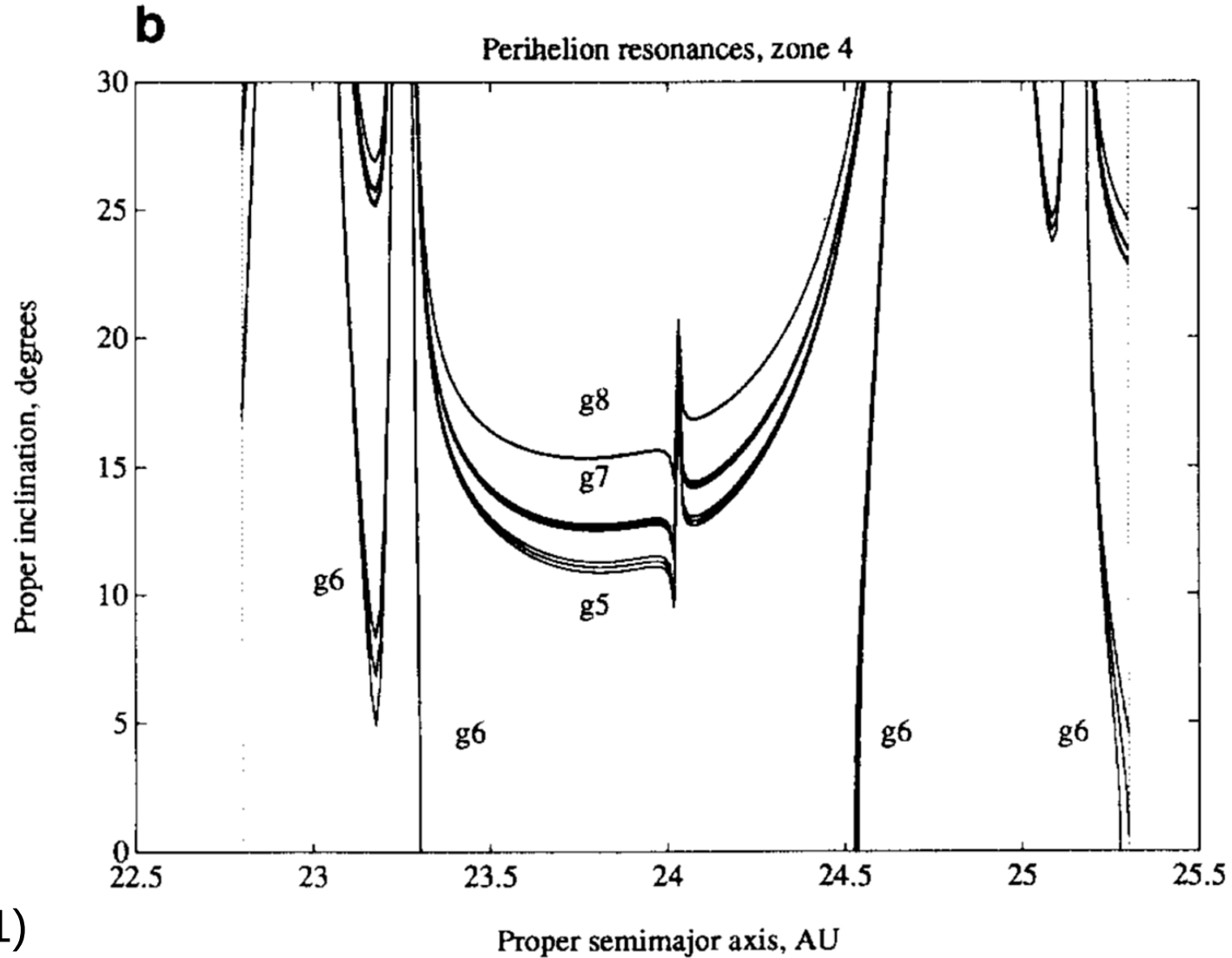
Examples - SR



- 1: $g - g_5 = 0$
- 2: $g - g_6 = 0$
- 3: $s - s_6 = 0$
- 4: $g + s - g_5 - s_6 = 0$
- 5: $g + s - g_6 - s_6 = 0$
- 9: $2g - g_5 - g_6 = 0$
- 10: $g - s - g_5 + s_6 = 0$
- 11: $g - s - g_6 + s_6 = 0$

Morbidelli & Henrard (1991)

Examples – SR



Knezevic,+ (1991)

2. Disturbing function

$E_k = \frac{1}{2} m v^2 \quad \text{tg } \vartheta_B = \frac{m_2}{m_1} = m_{21} \quad pV = nRT \quad \vec{\Psi} = \iint \vec{D} d\vec{S} = AD \quad H_\lambda = \frac{\Delta M_e}{\Delta \lambda}$
 $M_e = \sigma T^4 \quad \phi_e = \frac{L}{S} \int \frac{\Delta \varphi}{2\pi} = \frac{\Delta x}{\lambda} = \frac{x_2 - x_1}{\lambda} S_2 \quad v = c/\lambda \quad \Phi = NBS$
 $-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V\psi = E\psi \quad U_{ef} = \frac{U_m}{\epsilon - k \frac{v_0}{c}} \quad U = \frac{W_{AB}}{e} = \frac{|E_{PA} - E_{PB}|}{e} = |V_A - V_B| \quad T = \frac{4n_1 n_2}{(n_2 + n_1)^2}$
 $X_L = \frac{U_m}{I_m} = \omega L = 2\pi f L \quad F_g = \frac{m_1 m_2}{r^2} \quad R_m = \frac{c}{T} k = \pm \sqrt{\frac{2m}{\hbar^2} (E - V_0)}$
 $\vec{B} = \mu_0 \frac{NI}{2r} \quad v = \frac{wh}{2\pi r m_e} \quad \varphi_E = \frac{E_0}{k} \frac{q}{r} \quad \varphi = |V_A - V_B| \quad I = \frac{U_e}{R + R_i} \quad \omega = 2\pi f$
 $k = \frac{p}{\hbar} \quad m_0 = \frac{M_m}{N_A} = \frac{M_r \cdot 10^{-3}}{N_A} \quad m = N \cdot m_0 = \frac{Q}{e} \frac{M_m}{N_A} \quad E = \frac{E_c}{a} \int \sin(\omega t + \phi) dy$
 $\lambda = \frac{h}{p} \quad l_t = l_0(1 + d \Delta t) \quad \frac{\sin \alpha}{\sin \beta} = \frac{v_1}{v_2} = \frac{m_2}{m_1} \quad v = \frac{1}{\sqrt{\epsilon \mu}} = \frac{c}{\sqrt{\epsilon_r \mu_r}}$
 $\sqrt{2eU_m} \quad R = \rho \frac{l}{S} \quad E = mc^2 \quad \beta = \frac{\Delta I_c \phi_e}{\Delta t} = \frac{\Delta E}{\Delta t} \frac{m_1}{x} + \frac{m_2}{x'} = \frac{m_2 - m_1}{v}$
 $f_0 = \frac{1}{2\pi} \sqrt{\frac{g}{l}} \quad \psi(x) = \sqrt{2/L} \sin \frac{n\pi x}{L} \quad \vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B}) \quad E_k = \frac{h^2}{8mL^2} \frac{h^2}{\lambda^2} \quad \oint \vec{D} d\vec{S} = Q^*$
 $\oint \vec{B} d\vec{l} = \mu_0 \iint \vec{J} d\vec{S} \quad v_k = \sqrt{\frac{3kT}{m_0}} = \sqrt{\frac{3kTN_A}{M_m}} = \sqrt{\frac{3R_m T}{M_r \cdot 10^{-3}}}$
 $E = \hbar k^2 \quad 1 \text{ pc} = \frac{1 \text{ AU}}{r} \quad S = \frac{U}{I} \quad \psi_2 = U_e I t \quad \vec{F}_v = \int \frac{F_n}{R}$
 $\lambda = \frac{h v_2}{T} \quad F_h = Shp g \quad f_0 = \frac{1}{2\pi \sqrt{CL}} \quad S I_m^2 = U_m^2 \left[\frac{1}{R^2} + \left(\frac{1}{X_C} - \frac{1}{X_L} \right)^2 \right] \quad \lambda^* T = b$
 $\left(\frac{E_t}{E_0} \right)_{||} = \frac{2 \cos \vartheta_1 \cos \vartheta_2}{\cos(\vartheta_1 - \vartheta_2) \sin(\vartheta_1 + \vartheta_2)} \quad \int \vec{E} d\vec{l} = - \iint \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} \quad \rho = \frac{E}{c} = \frac{hf}{c} = \frac{h}{\lambda}$
 $E_y = E_0 \sin(kx - \omega t) \quad R = R_0 \sqrt{A} \quad S = \frac{1}{A} \frac{dW}{dt} \quad \oint \vec{H} d\vec{l} = \iint (\vec{J} + \frac{\partial \vec{D}}{\partial t}) \cdot d\vec{S} \quad \varphi = mc \Delta t \quad F_g = \frac{M_0 M_2}{r^2}$
 $W = F \cdot s \cdot \cos \alpha \quad L = 10 \log \frac{I}{I_0} \quad \mu = U_m \sin \omega(t - T) = U_m \sin 2\pi \left(\frac{t}{T} - \frac{x}{\lambda} \right)$
 $\oint \vec{B} d\vec{l} = \mu_0 \sum I; \quad \rho = \frac{\vec{F}}{\Delta S} = \frac{m \Delta \vec{v}}{\Delta S \Delta t} \quad \Delta \psi = \frac{2\pi \Delta x}{\lambda} = \frac{2\pi d \sin \vartheta}{\lambda} = \frac{2\pi dy}{xL}$
 $R = \frac{(n-1)^2 + g^2}{(n+1)^2 + g^2} \quad f' = \frac{n_a \cdot n_b}{(n-1)(n_0 - n_a)} \quad \nabla \times \left(-\frac{\partial \vec{B}}{\partial t} \right) = \frac{\partial}{\partial t} (\text{rot } \vec{B}) = -\mu_0 \frac{\partial}{\partial t} \left(\frac{\partial \vec{B}}{\partial t} \right) = \epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2}$

Disturbing function – basics

- Basis for analytical approximations for $N \geq 3$ bodies
- Gradient of perturbing potential
- Equations of motion in heliocentric Cartesian coordinates

$$\ddot{\mathbf{r}}_i = -G(M + m_i) \frac{\mathbf{r}_i}{\|\mathbf{r}_i\|^3} + \sum_{n=1}^N Gm_n \left(\frac{\mathbf{r}_n - \mathbf{r}_i}{\|\mathbf{r}_n - \mathbf{r}_i\|^3} - \frac{\mathbf{r}_n}{\|\mathbf{r}_n\|^3} \right)$$

Disturbing function – variables

- Express disturbing function in orbital elements
- Definition (Stiefel & Scheifele, 1971)
an **orbital element** in the 2-body problem
is a **linear function of time**: $f(t) = a + b t$
 - $(a, e, i, \omega, \Omega) = \text{const.}$
 - $M = n (t - t_0) = M_0 + n t \dots$ only time dependent function
- Slow **non-linear** time-variation for $N > 2$ body problem

Disturbing function – series expansions

Development in different coordinate systems:

- Heliocentric coord.
 - Murray & Dermott (1999), Ellis & Murray (2000)
- Mixed barycentric–heliocentric coord.
 - Laskar & Robutel (1995)
- Jacobi coord.
 - Mardling (2013)

Disturbing function – application

- Lagrange planetary equations
- Express time variation of orbital elements by:

$$\frac{da}{dt} = \frac{2}{na} \frac{\partial \mathcal{R}}{\partial \varepsilon}$$

$$\frac{de}{dt} = -\frac{\sqrt{1-e^2}}{na^2e} \left(1 - \sqrt{1-e^2}\right) \frac{\partial \mathcal{R}}{\partial \varepsilon} - \frac{\sqrt{1-e^2}}{na^2e} \frac{\partial \mathcal{R}}{\partial \varpi}$$

$$\frac{di}{dt} = -\frac{\tan(i/2)}{na^2\sqrt{1-e^2}} \left(\frac{\partial \mathcal{R}}{\partial \varpi} + \frac{\partial \mathcal{R}}{\partial \varepsilon}\right) - \frac{1}{na^2\sqrt{1-e^2} \sin i} \frac{\partial \mathcal{R}}{\partial \Omega}$$

$$\frac{d\varpi}{dt} = \frac{\sqrt{1-e^2}}{na^2e} \frac{\partial \mathcal{R}}{\partial e} + \frac{\tan(i/2)}{na^2\sqrt{1-e^2}} \frac{\partial \mathcal{R}}{\partial i}$$

$$\frac{d\Omega}{dt} = \frac{1}{na^2\sqrt{1-e^2} \sin i} \frac{\partial \mathcal{R}}{\partial i}$$

$$\frac{d\varepsilon}{dt} = -\frac{2}{na} \frac{\partial \mathcal{R}}{\partial a} + \frac{\sqrt{1-e^2}(1 - \sqrt{1-e^2})}{na^2e} \frac{\partial \mathcal{R}}{\partial e} + \frac{\tan(i/2)}{na^2\sqrt{1-e^2}} \frac{\partial \mathcal{R}}{\partial i}$$

3. Mean-motion resonance (MMR)

MMR – terms in disturbing function

Table B.1. Zeroth-order arguments: direct part.

ID	Cosine Argument	Term
4D0.2	$j\lambda' - j\lambda + \varpi' - \varpi$	$ee' f_{10} + e^3 e' f_{11} + ee'^3 f_{12}$ $+ ee'(s^2 + s'^2) f_{13}$
4D0.3	$j\lambda' - j\lambda + \Omega' - \Omega$	$ss' f_{14} + ss'(e^2 + e'^2) f_{15}$ $+ ss'(s^2 + s'^2) f_{16}$
4D0.4	$j\lambda' - j\lambda + 2\varpi' - 2\varpi$	$e^2 e'^2 f_{17}$
4D0.5	$j\lambda' - j\lambda + 2\varpi - 2\Omega$	$e^2 s^2 f_{18}$
4D0.6	$j\lambda' - j\lambda + \varpi' + \varpi - 2\Omega$	$ee' s^2 f_{19}$
4D0.7	$j\lambda' - j\lambda + 2\varpi' - 2\Omega$	$e'^2 s^2 f_{20}$
4D0.8	$j\lambda' - j\lambda + 2\varpi - \Omega' - \Omega$	$e^2 ss' f_{21}$
4D0.9	$j\lambda' - j\lambda + \varpi' - \varpi - \Omega' + \Omega$	$ee' ss' f_{22}$
4D0.10	$j\lambda' - j\lambda + \varpi' - \varpi + \Omega' - \Omega$	$ee' ss' f_{23}$
4D0.11	$j\lambda' - j\lambda + \varpi' + \varpi - \Omega' - \Omega$	$ee' ss' f_{24}$
4D0.12	$j\lambda' - j\lambda + 2\varpi' - \Omega' - \Omega$	$e'^2 ss' f_{25}$
4D0.13	$j\lambda' - j\lambda + 2\varpi - 2\Omega'$	$e^2 s'^2 f_{18}$

Murray & Dermott (1999)

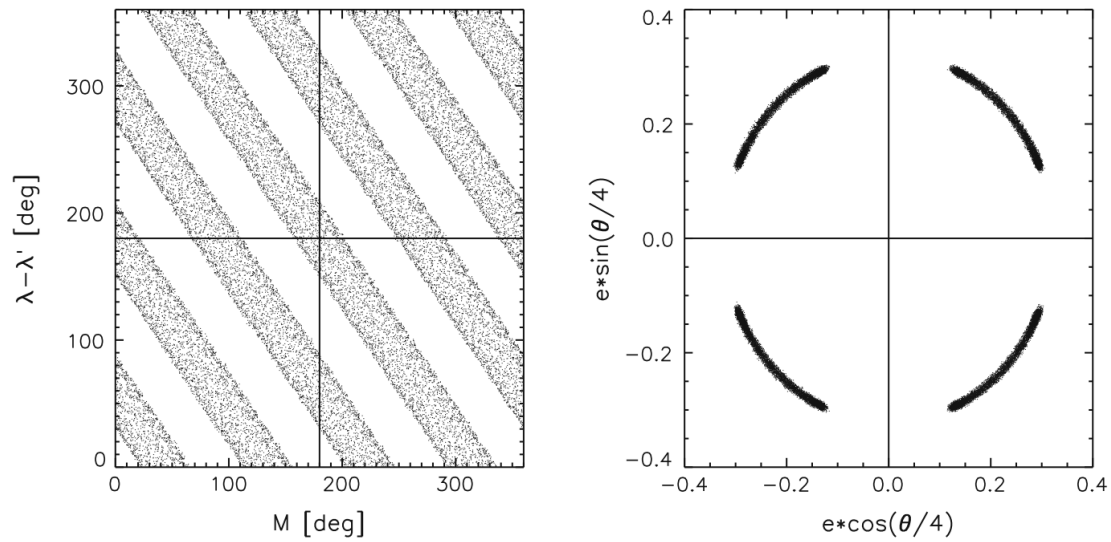
MMR – theoretical concepts

- MMR → orbital frequencies
- Critical angle of MMR
- Small divisor for $j_1 n + j_2 n' \approx 0$
- Resonance location a_{res} from 3rd Kepler law

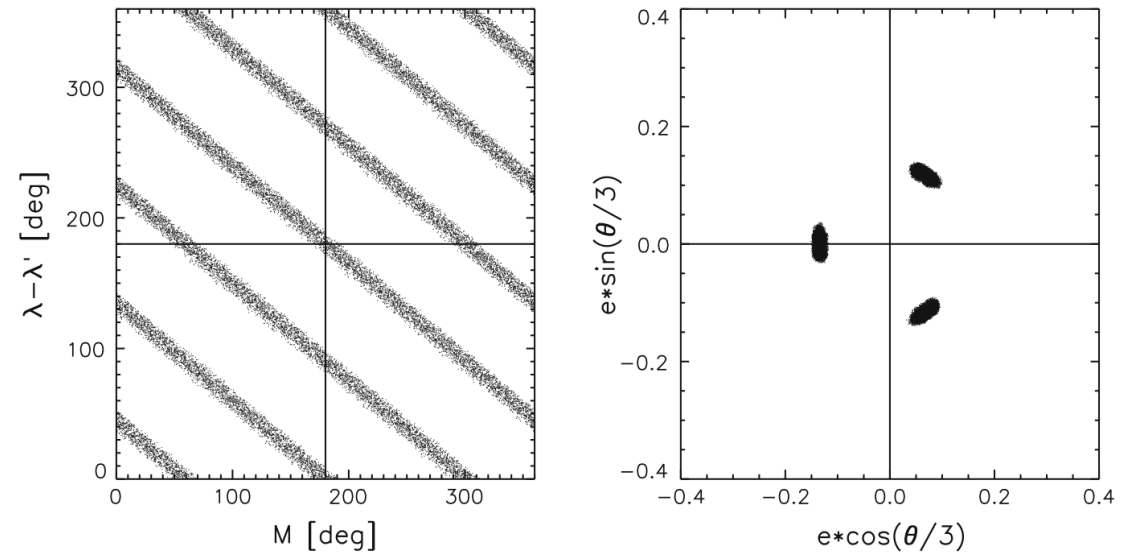
$$a_{\text{res}} = a' \left(\frac{n'}{n} \right)^{2/3} \left(\frac{M + m}{M + m'} \right)^{1/3}$$

MMR – visualization

7:3 MMR

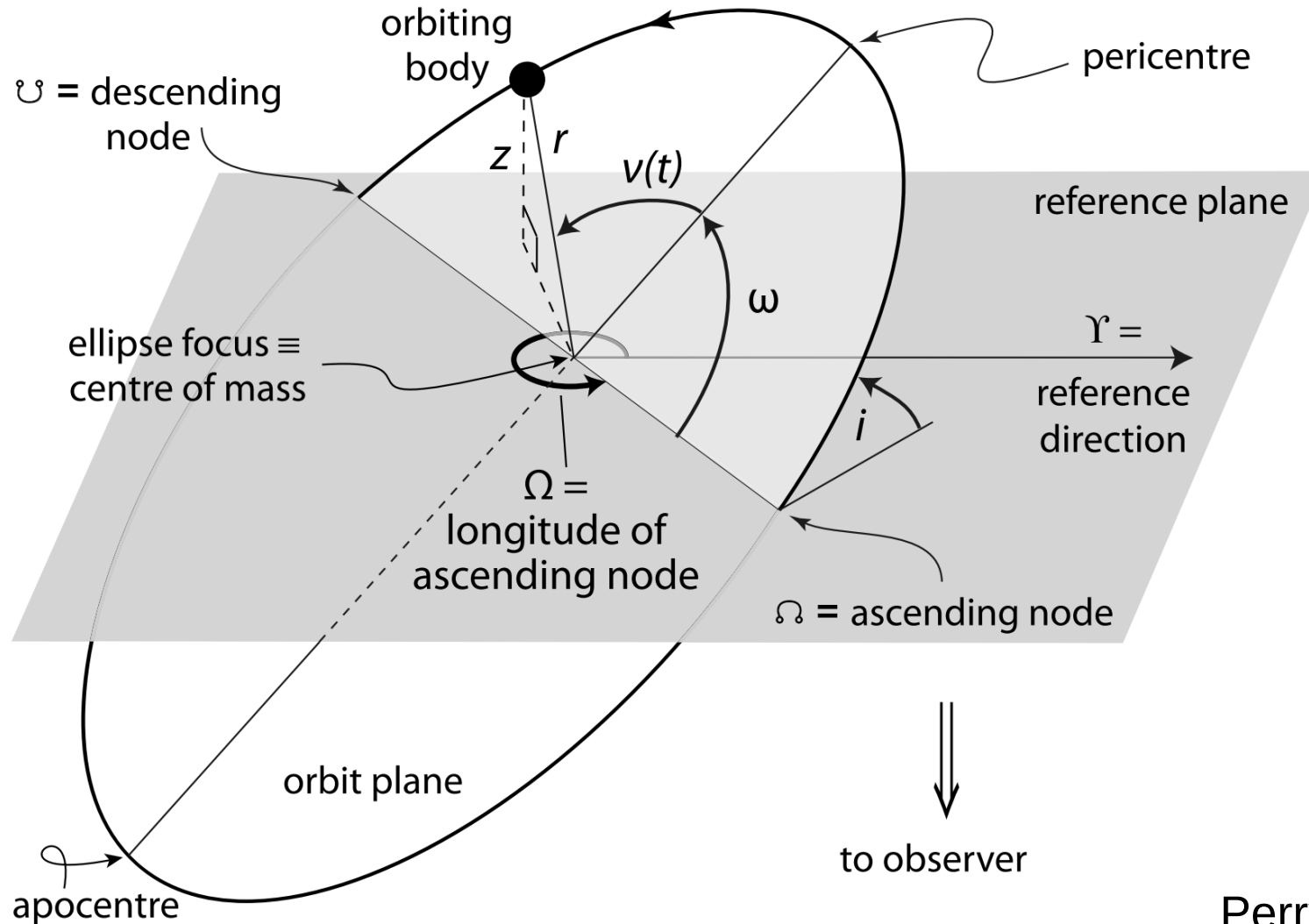


7:4 MMR

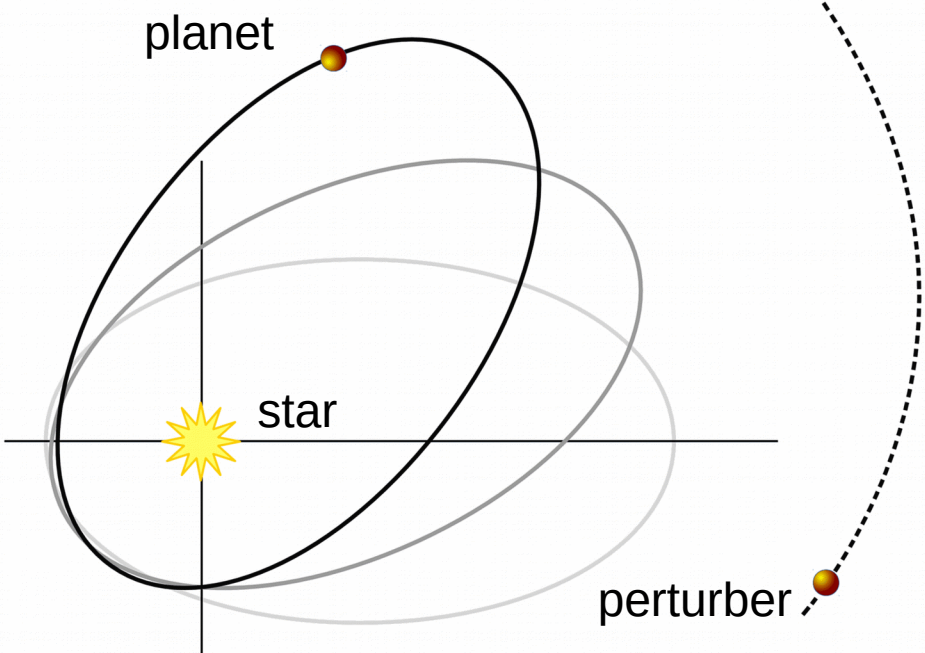
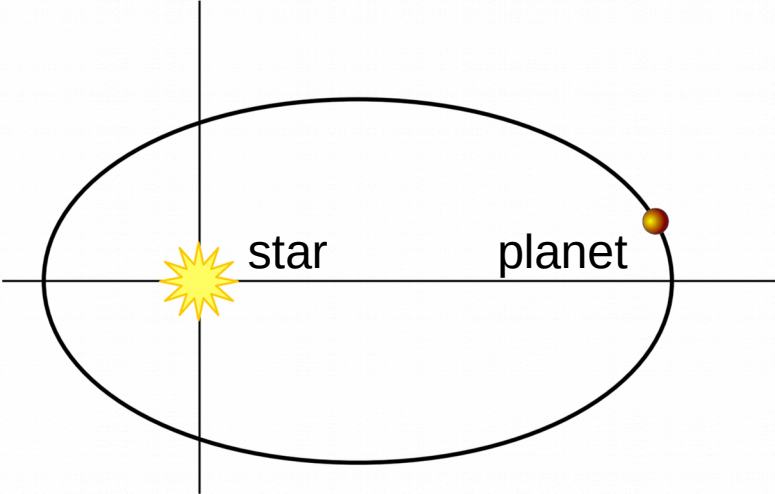


4. Secular resonance (SR)

Secular resonances



Secular resonances



SR – theoretical concepts

- SR → orbital precession frequencies
- Line of apsides, line of nodes
- Time-scale $T_{\text{sec}} \gg T_{\text{rev}}$

SR – disturbing function

- Averaging principle
- Remove short period contributions
- Eliminate “fast” frequencies $\lambda \sim M$
- Averaged (secular) disturbing function

$$\mathcal{R}(a, e, i, \omega, \Omega, \lambda) \longmapsto \langle \mathcal{R} \rangle(-, e, i, \omega, \Omega, -) = \frac{1}{2\pi} \int_0^{2\pi} \mathcal{R}(a, e, i, \omega, \Omega, \lambda) d\lambda$$

SR – secular variables

- Laplace-Lagrange variables
- Decoupling of eccentricity / inclination (to lowest order)

$$h = e \sin(\omega + \Omega)$$

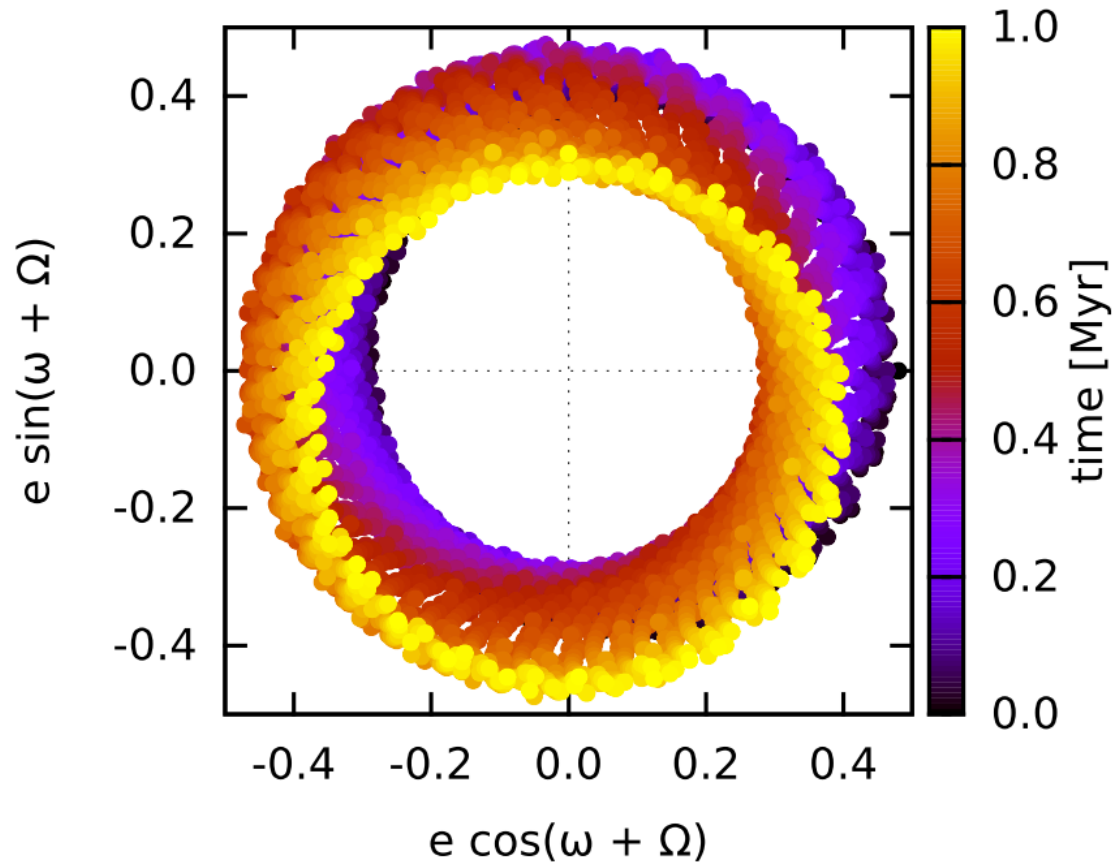
$$k = e \cos(\omega + \Omega)$$

$$p = \sin(i/2) \sin \Omega$$

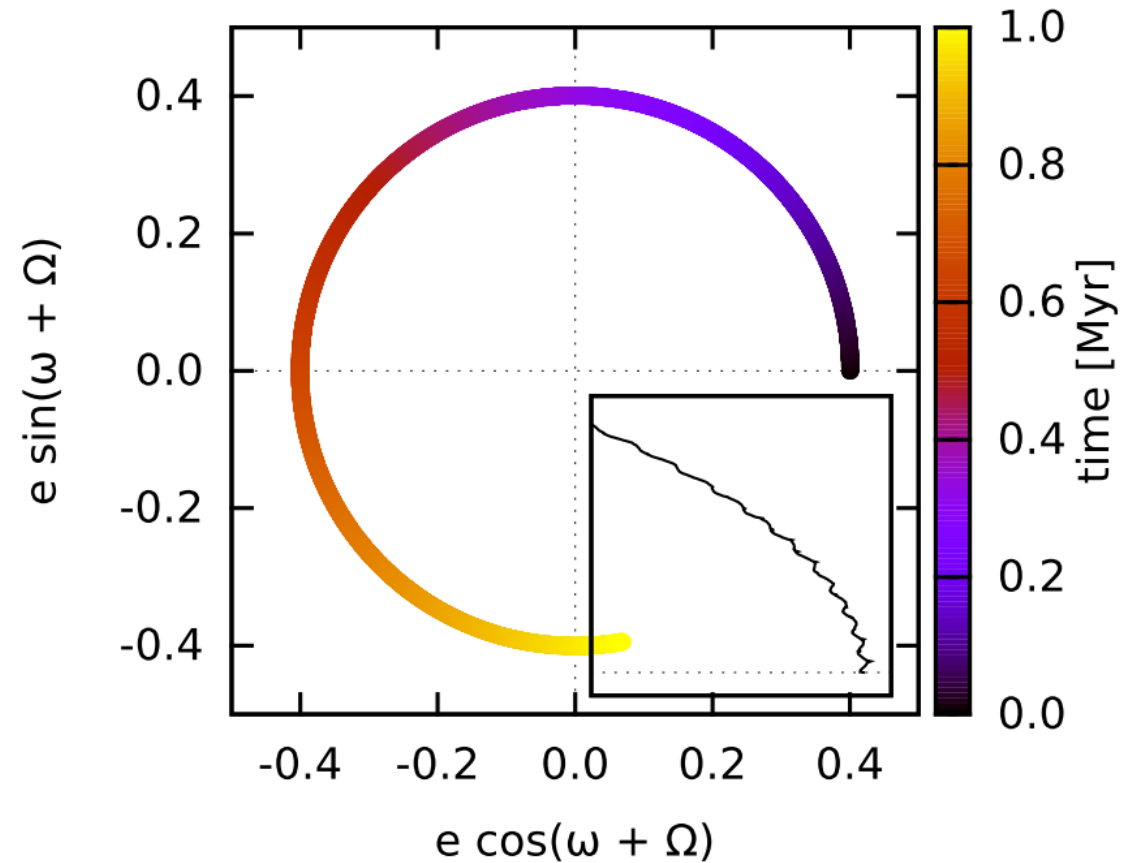
$$q = \sin(i/2) \cos \Omega$$

SR – orbital precession

giant planet



secondary



SR – solutions

- Equations of motion in variables (h,k) – system of linear differential equations
- Secular eigenfrequencies (eigenvalues) g_i
- Laplace coefficients $b_n^{(k)}(\alpha)$

$$\dot{\mathbf{h}} = \mathbf{A}\mathbf{k}$$

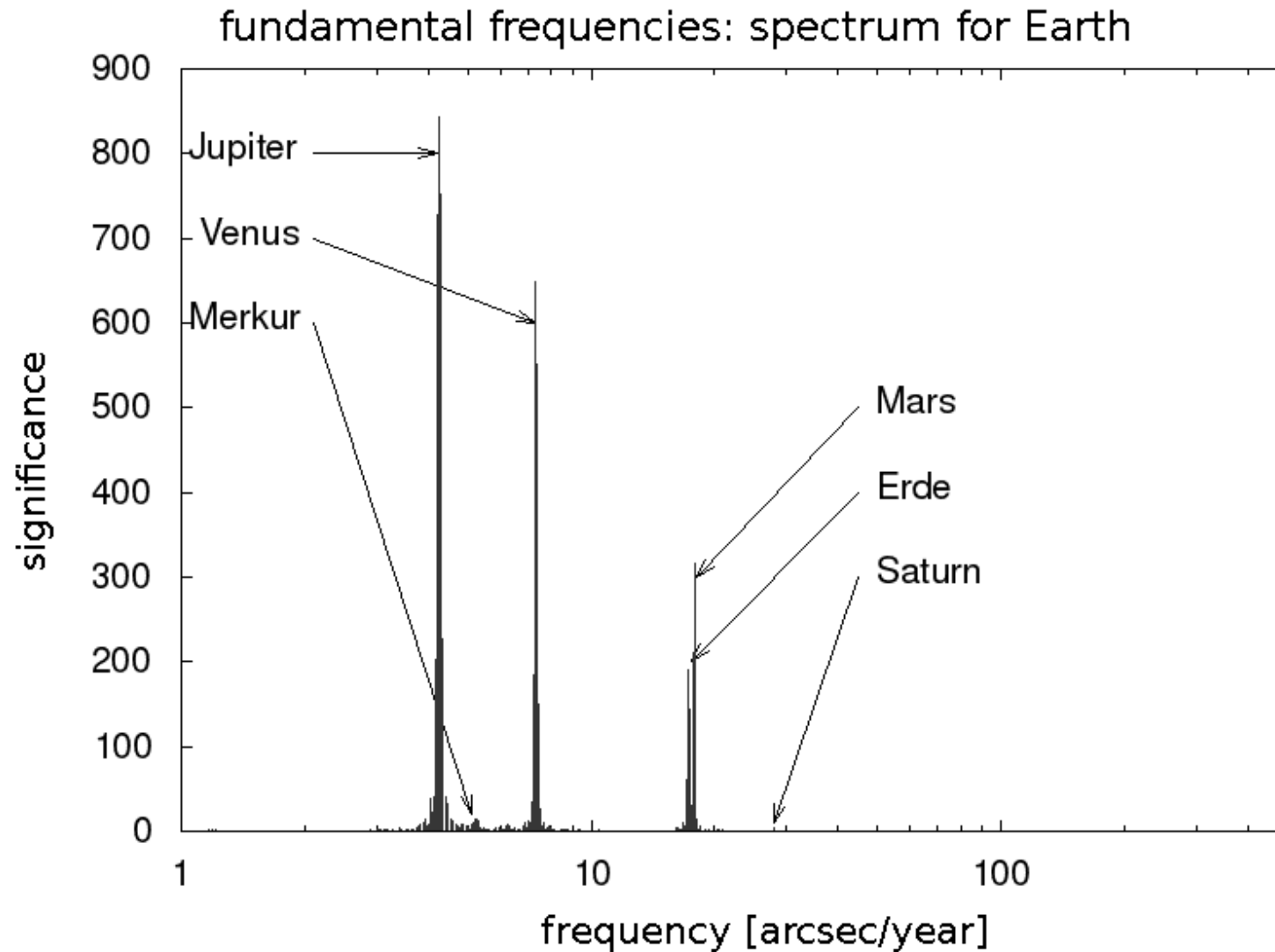
$$\dot{\mathbf{k}} = -\mathbf{A}\mathbf{h}$$

$$A_{j,j} = +\frac{1}{4}n_j \sum_{k=1}^N \frac{m_k}{M + m_j} \alpha_{j,k}^2 b_{3/2}^{(1)}(\alpha_{j,k})$$

$$A_{j,k} = -\frac{1}{4}n_j \frac{m_k}{M + m_j} \alpha_{j,k}^2 b_{3/2}^{(2)}(\alpha_{j,k})$$

$$\det(\mathbf{A} - g\mathbf{1}) = 0$$

SR – frequencies



SR – test particle

- Disturbing function for a TP with N massive perturbers
- **Proper frequency** g of TP
- General solution for TP in (h,k) variables
- **Small divisor** for $g - g_i \approx 0$
- Proper (free) + forced eccentricity / inclination

$$\mathcal{R} = n a^2 \left[\frac{1}{2} g (h^2 + k^2) + \sum_{j=1}^N A_j (h h_j + k k_j) \right]$$

$$g = \frac{1}{4} n \sum_{j=1}^N \frac{m_j}{M} \alpha_j^2 b_{3/2}^{(1)}(\alpha_j)$$

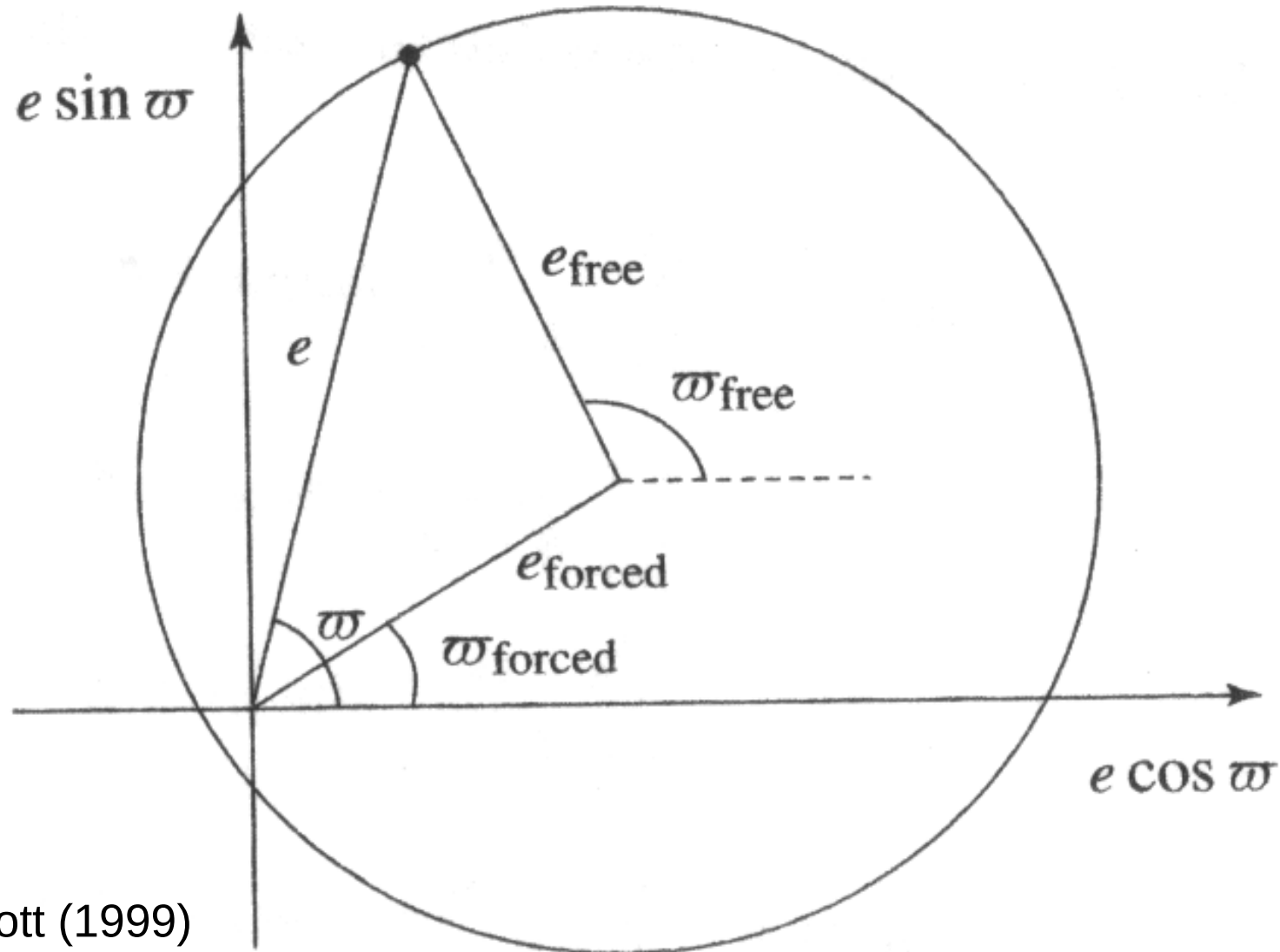
$$\begin{aligned} h(t) &= e_{\text{free}} \sin(gt + \varphi) - \sum_{i=1}^N \frac{\nu_i}{g - g_i} \sin(g_i t + \varphi_i) \\ &= h_{\text{free}}(t) + h_0(t) \end{aligned}$$

$$\begin{aligned} k(t) &= e_{\text{free}} \cos(gt + \varphi) - \sum_{i=1}^N \frac{\nu_i}{g - g_i} \cos(g_i t + \varphi_i) \\ &= k_{\text{free}}(t) + k_0(t) \end{aligned}$$

$$e_{\text{forced}} = \sqrt{h_0^2 + k_0^2}$$

$$i_{\text{forced}} = \sqrt{p_0^2 + q_0^2}$$

SR – free / forced eccentricity



Murray & Dermott (1999)

Summary

- **Concepts:**
 - What is a resonance?
 - Which kinds of resonances exist?
- **Disturbing function:**
 - What is it representing?
 - Why is it useful?
- **Mean-motion resonances:**
 - What causes MMR?
- **Secular resonances:**
 - What causes SR?

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