On the dynamical stability of the solar system

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Introduction

- chaotic behavior of solar system (Laskar 1989)
- 4 terrestial planets: 5 mio. yrs Lyapunov-time (Laskar 1989)
- mass ratios of the planets to the sun are much larger than those required from KAM
- ♦ Is the solar system dynamically stable?

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Introduction

- long-term numerical integration of the full solar system over 20 Gyrs
- dynamically allowed evolutions in which the planetary orbits become unstable
- effects of general relativity on the dynamical stability
- dynamical lifetime of Uranus

chaotic motion - Lyapunov time $\gamma = \lim_{N \land \infty} \sum_{k=1}^{N} \ln (s_k / s_0) / (N\Delta t)$

s₀ ... separation vector (150 m, radially outward) Δt... 10.000 yr N... 100 $1/\gamma$... Lyapunov time $s/s_0 = e^{(\gamma \Delta t)}$

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Table 1: Lyapunov exponents and times for the Solar System

Planet	Lyapunov	exp. (yrs 10 ⁻¹)	Lyapunov time	(yrs)
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Mercury	7.32029 10-7	$1.36607 \ 10^{6}$
Venus	1.38561 10-7	7.21703 106
Earth	2.07484 10-7	4.81964 106
Mars	2.22353 10-7	4.49736 106
Jupiter	1.19528 10-7	8.36623 106
Saturn	1.56875 10-7	6.37452 106
Uranus	1.33793 10-7	7.47423 106
Neptune	1.49602 10-7	6.68440 106

chaotic motion - Lyapunov time

• Lyapunov time - some 10⁶ yrs

Why long-term integrations?

They give probalistic evaluations of the solar system's future behavior.

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direct long-term integration

- start: current configuration
- timespan: over 20 Gyrs
- timestep: 8 days
- conservation of total energy: $\Delta E/E < 10^{-7}$
- conservation of total angular momentum: $\Delta L/L < 10^{-9}$

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Eccentricity of earth as function of time

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Eccentricity of mercury as function of time

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direct long-term integration

solar system seems to be stable over its
lifetime?

but: its only one possible trajectory! (Laskar 1994)

 better: any timescale for occurring instabilities might be long

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The Laskar - method



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The Laskar - method

numerical induced chaos - limiting factor: accurately resolving Mercury's orbit(Yoshida, 1993) conservation of energy and ang.momentum
 ◊ criteria for this work: $\Delta E/E \sim 10^{-8}$ and $\Delta L/L \sim 10^{-10}$ ♦ variation of timestep form 3 to 1.2 days to correct the violation of these criterias

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• 2 experiments: 150m and 15 m

Step	Time interval	Endpoint
1	0-500	0.2907
2	500-797	0.4391
3	797-862	0.8257

Step	Time interval	Endpoint
_1	0-500	0.2907
2	500-994	0.4139
3	994-1207	0.4874
4	1207-1261	0.9751

Ollision with Venus

♦ Collision with Sun

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Fig. 10. The minimal distance of approach during a series of close encounters between Mercury and Venus as a function of time. The collision takes place at t ~ 861.455Myr, when dmin = $5.5561 \times 10^{-5} \text{ AU} < r_venus + r_mercury = 5.6762 \times 10^{-5} \text{ AU}.$

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Evolution of Mars' eccentricity in the 150 m experiment at t = 822 Myrs - also Mars' semimajor axis increased and after it reached a distance > 100 AU, Mars was assumed to be ejected fromsolar system

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- Mercury shows different ways of evolution
- Mercury tends to switch from regular to irregular motion
- Exists any association with timesteps or integral algorithm?
- Reintegration of a 22 Myr time interval where Mercury became unstabel

<u>150 m</u>

symplectic algorithm0.5 days time step

• 22 Myrs

Bulirsch-Stoer algorithm 0.5 days time step

<u>15 m</u>

symplectic algorithm0.5 days time step

• 20 Myrs

Bulirsch-Stoer algorithm 0.5 days time step

almost identical increase of eccentricity as in the primary solution

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Starting at 778 Myrs for the 150 m:

- adding perturbations of 15 m and 150 m in four directions and integrating over 22 Myrs
- ◊ no changes

Starting at 1190 Myrs for the 15 m:

 adding perturbations of 15 m and 150 m in four directions and integrating over 20 Myrs
ho changes

- until 780 Myrs Mercury's eccentricity varies within a narrow and well-defined range
- shortly thereafter the eccentricity increases and leads to collisions,.....
- ◊ secular resonances

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stable motion versus unstable motion

$$<\!\!R_i^{sec}\!\!> = n_i a_i^2 \{ 0.5 A_{ii} e_i^2 + 0.5 B_{ii} I_i^2 + \sum [A_{ij} e_i e_j \cos(\overline{\omega}_i - \overline{\omega}_j) + B_{ij} I_i I_j \cos(\Omega_i - \Omega_j)] \}$$

 $R_{i} = \sum \left[(G m_{j}) / (|r_{j} - r_{i}|) - G m_{j} (r_{j} \times r_{i}) / (r_{j}^{3}) \right]$

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The dominant frequencies of the secular disturbing function can be identified by Fourier-analyzing the numerically computed time-series for Mercury's full disturbing function.

$$g_i = \langle \varpi_i \rangle$$

stable motionunstable motion $(g_1 - g_5) \sim 0.9389$ " yr⁻¹ $(g_1 - g_5) \sim 0.0538$ " yr⁻¹

(Laskar 1990)

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secular resonances - stable motion



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secular resonances - unstable motion



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Classical Laplace-Lagrange secular solution for Mercury's ecc. vector:

$$h_1 = e_1 \sin \varpi_1$$
$$k_1 = e_1 \cos \varpi_1$$

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$$h_1 = e_{1f} \sin (g_1 t + \beta_1) - \sum [(v_j / (g_1 - g_j)) \sin (g_j t + \beta_j)]$$

$$k_1 = e_{1f} \cos (g_1 t + \beta_1) - \sum [(v_j / (g_1 - g_j)) \cos (g_j t + \beta_j)]$$

small divisors indicate large influence !!

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classical value of Mercury's proper frequ.: $g_1 = 5.4058$ " during Mercury's evolution of ecc.: $g_1 = 4.9273$ " $g_i = 4.24354$ " ◊ linear secular resonances ◊ large variations in the ecc. vector

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 $de_{15}/dt = A_{15} e_5 (1-e^2)^{0.5} sin(\varpi_1 - \varpi_5)$

 $g_1 \Rightarrow \langle \varpi_1 \rangle \text{ and } g_5 \Rightarrow \langle \varpi_5 \rangle$

 ◊ (g₁ - g₅) secular resonance responsible for the climb in Mercury's eccentricity.

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Influence of Venus:

$g_2 \not\Rightarrow \langle \varpi_2 \rangle$ (Laskar, 1990)

 $\diamond e_{12}$ constructed as before

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relativistic influence

- adding 0.43" yr⁻¹
- $g_1 g_5$ gets about 40 % higher
- Mercury enters later in the linear secular resonance with Jupiter
 relativistic influence stabilizes Mercury 's orbit

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Uranus dynamical lifetime

- Laskar experiment for Uranus
- sun (+4 terrestrial planets) and the 4 gasgiants
- Jupiter was repositioned 1500m
- Integration intervals: 5 Gyrs

◊ Uranus ecc. never exceeded 0,078 !!

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