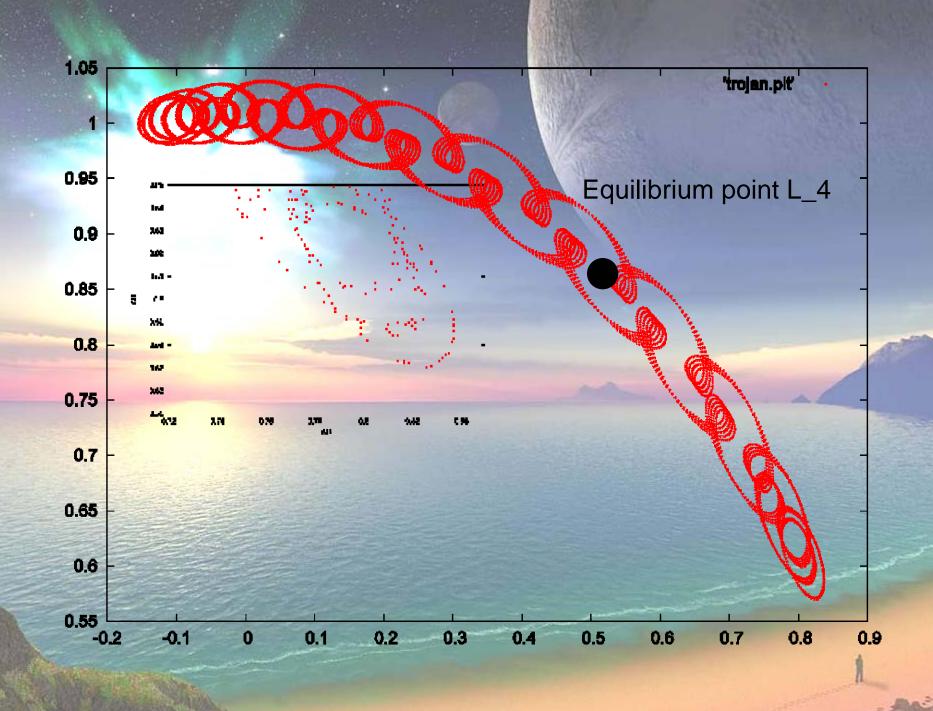


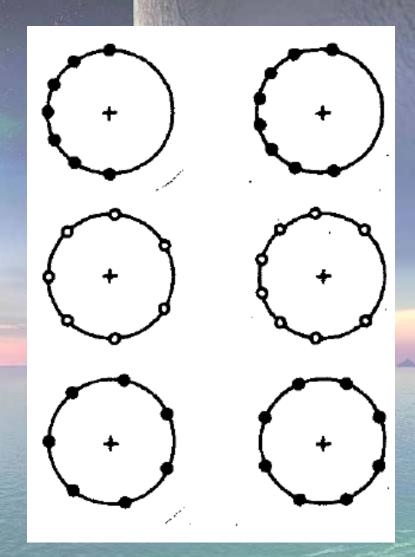
5th Austrio-Hungarian Workshop

April, 9th - April, 10th Vienna Observatory

The five Lagrange points SUN + JUPITER L_2 L_3



What happens when we have not only 2 equilibrium points but several- a whole ring of celestial bodies around a central body?



The dynamics of coorbital satellite systems

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Summary. The dynamical behavior of N coorbital satellites, moving with the same average mean motion around a primary has been studied both analytically and in terms of numerical integrations for $2 \le N \le 9$ satellites. Simplified dynamical equations have been used to determine the different stationary configurations and their local stability against infinitesimal perturbations. The motion is reduced to angular separations between satellites and is accurate to the first order in satellite to primary mass ratio. The ring of equally spaced identical satellites is found to be locally unstable for $N \le 6$, while for $2 \le N \le 8$ there exists another, stable compact stationary configuration, with separations $\leq 60^{\circ}$ between adjacent satellites. For $N \geq 7$ the equally spaced configuration becomes locally stable, and for $N \geq 9$ it is the only stationary configuration. Exact integrations

2. Simplified dynamical equations

The equations of motion for N mutually gravitating satellites, moving around a central mass in coplanar orbits, can be written in the form (see e.g. Yoder et al., 1983; Brown and Shook, 1964),

$$\frac{d}{dt}\left(r_i^2 \frac{d\theta_i}{dt}\right) = \sum_{\substack{j=1\\j\neq i}}^N \frac{\partial F_{ij}}{\partial \theta_i},\tag{1}$$

$$\frac{d^2r_i}{dt^2} - r_i \left(\frac{d\theta_i}{dt}\right)^2 = -\frac{\gamma M}{r_i^2} + \sum_{\substack{j=1\\j\neq i}}^N \frac{\partial F_{ij}}{\partial r_i},\tag{2}$$

where

$$F_{ij} = \gamma m_j \left\{ \frac{1}{\Delta_{ij}} - \frac{r_i}{r_j^2} \cos(\theta_i - \theta_j) \right\}, \tag{3}$$

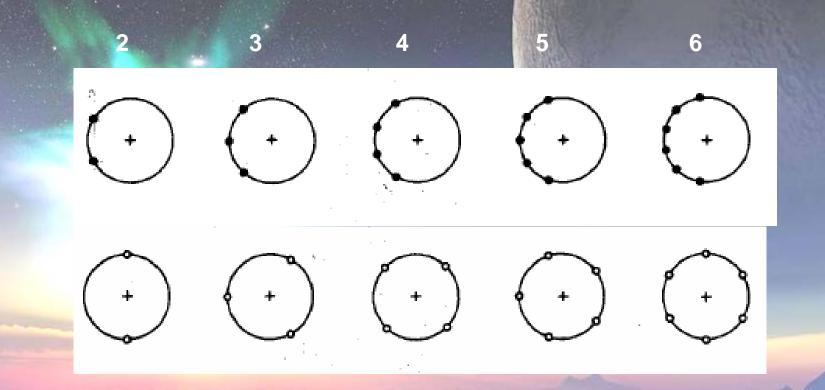
$$\Delta_{ij}^{2} = r_{i}^{2} + r_{j}^{2} - 2r_{i}r_{j}\cos(\theta_{i} - \theta_{j}). \tag{4}$$

The polar coordinates (r, θ) are referred to the center of mass of the primary, which gives rise to the indirect terms $\sim \cos(\theta_i - \theta_j)$ in the forcing functions F_{ij} . Satellite masses are denoted by m_i , while M stands for the central mass and γ for the gravitational constant.

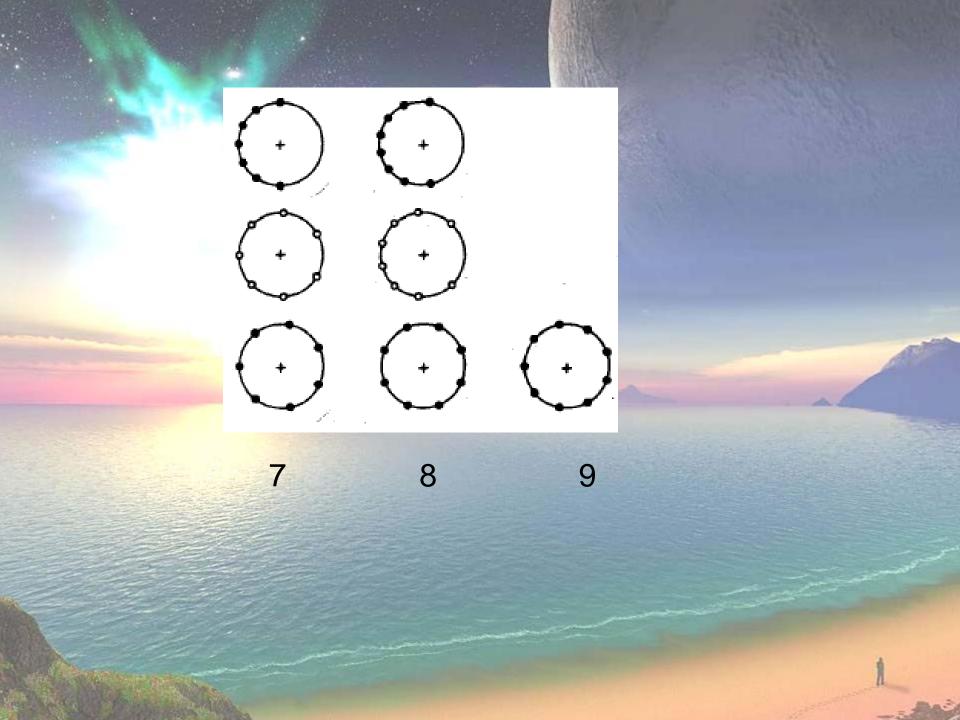
For N > 2 the situation is more complex, due to the large number of dependent variables: the energy condition is not sufficent to determine all the difference angles. However, for N = 3 some progress can be made by plotting the mutual gravitational potential as a function of angular separations, $\phi_i = \theta_{i+1} - \theta_i$. For convenience, assume $\theta_1 < \theta_2 < \theta_3$, and define new variables,

$$\alpha = (\phi_1 + \phi_2)/2 = (\theta_3 - \theta_1)/2; \qquad 0^{\circ} < \alpha < 180^{\circ},
\beta = (\phi_1 - \phi_2)/2 = \theta_2 - (\theta_3 + \theta_1)/2; \qquad -\alpha < \beta < \alpha,$$
(12)

so that α corresponds to the mean interparticle separation, while β describes the deviation of the middle particle from the equidistant position. Figure 1a shows the potential surface as a func-



Full cercles -stable orbits



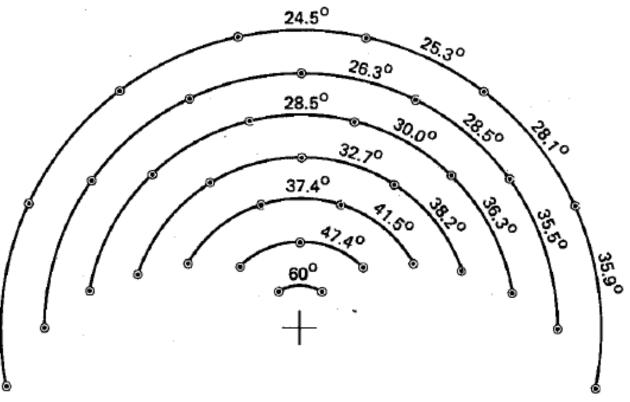
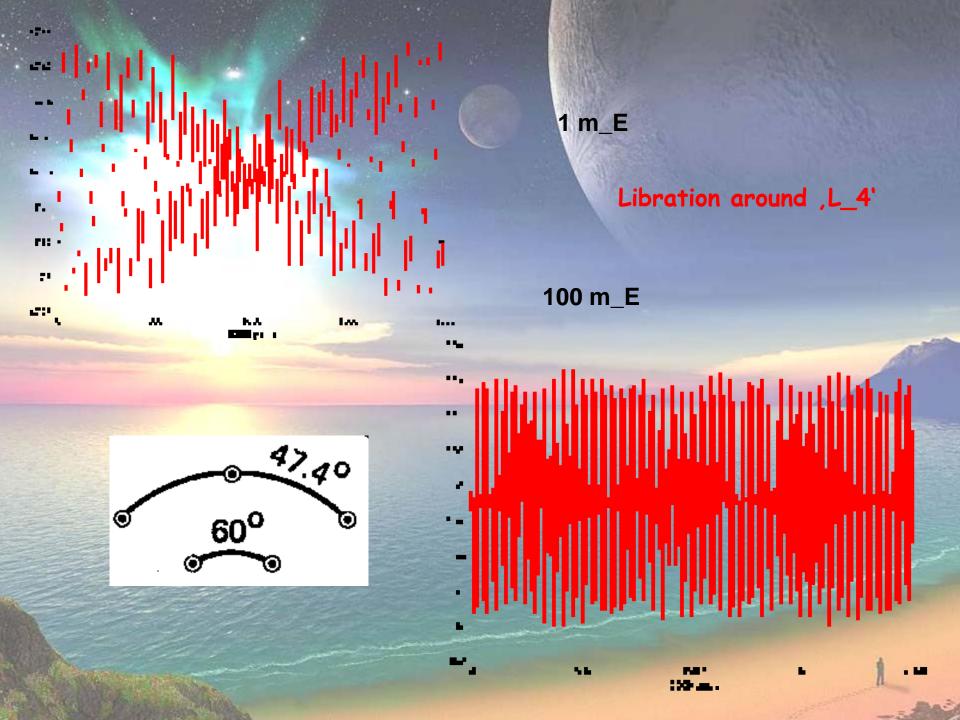
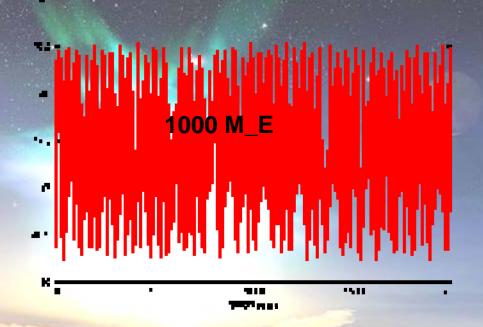
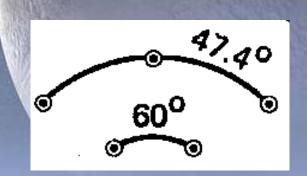


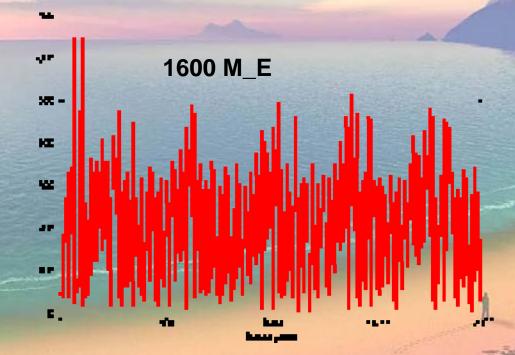
Fig. 3. The angular separations for the stable compact configuration (Type I), for $2 \le N \le 8$ satellites

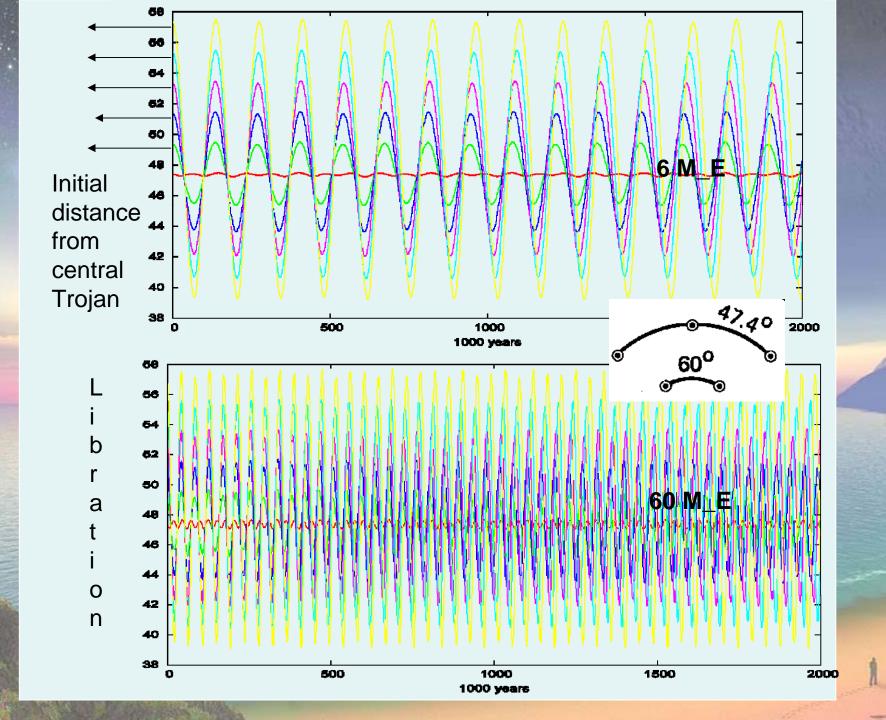


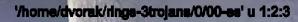


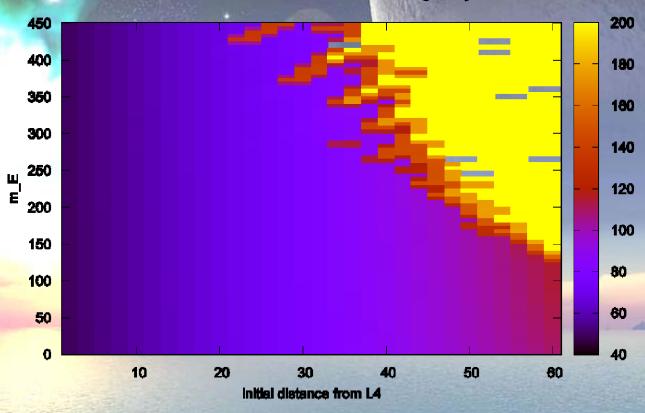


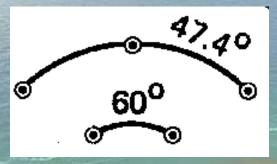


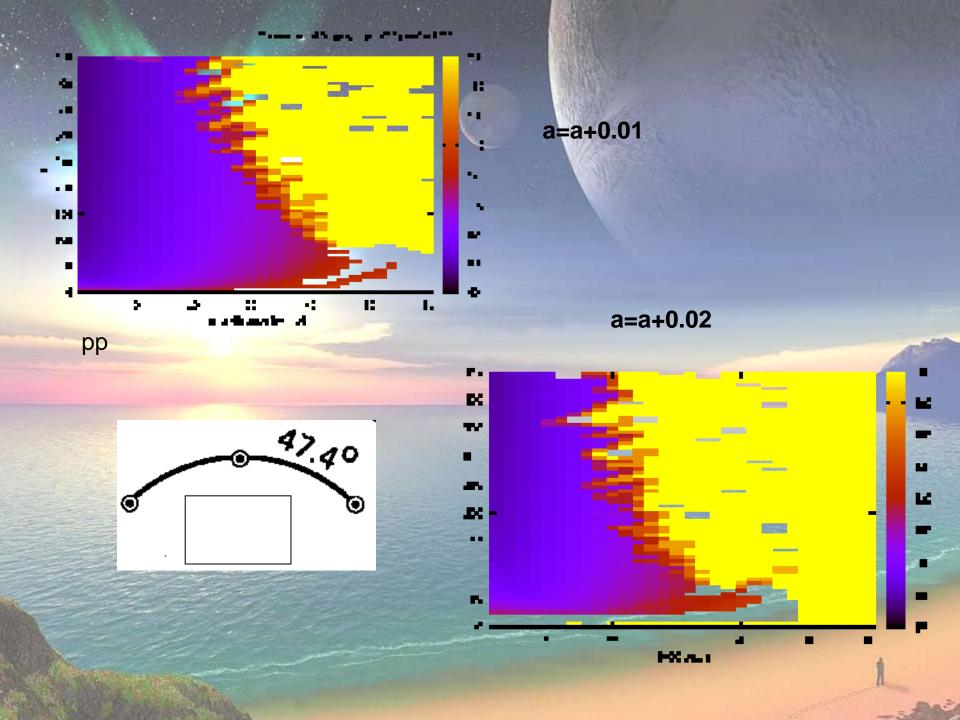


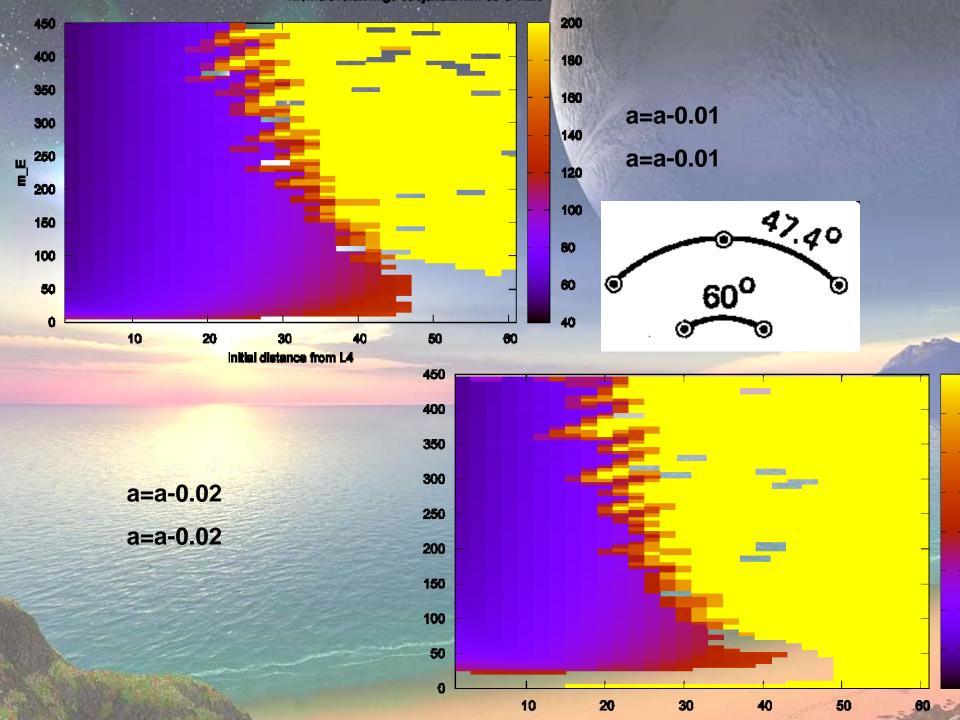


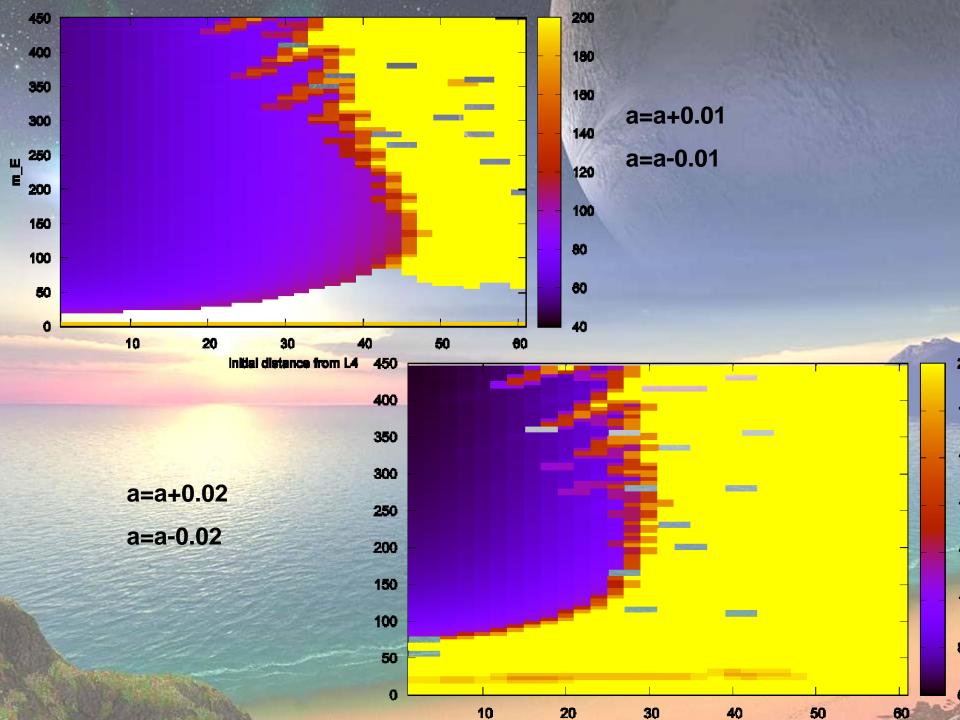


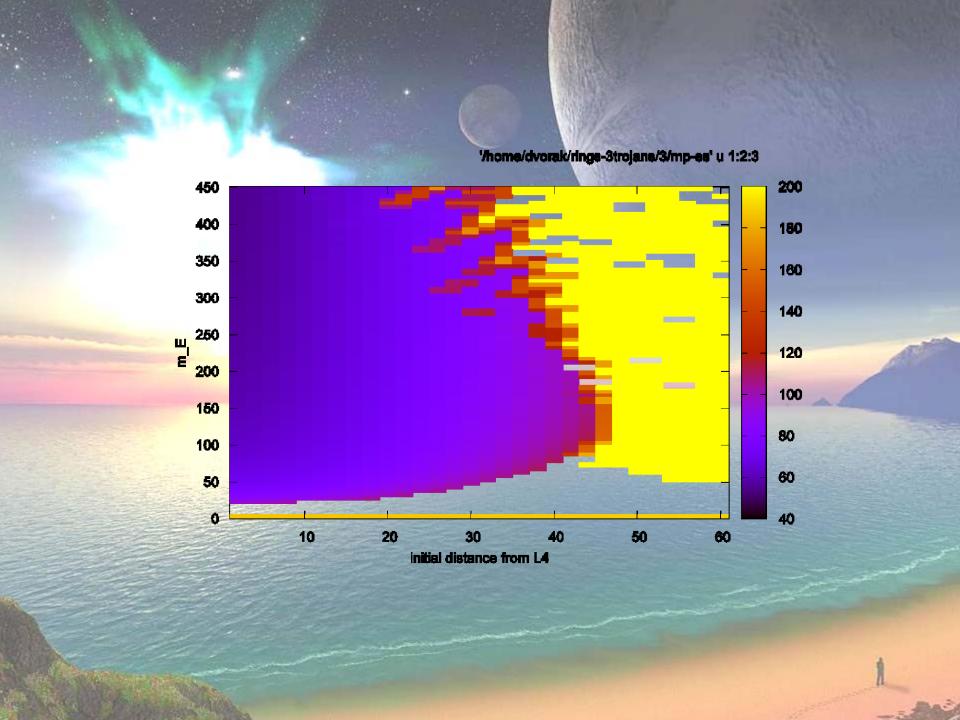










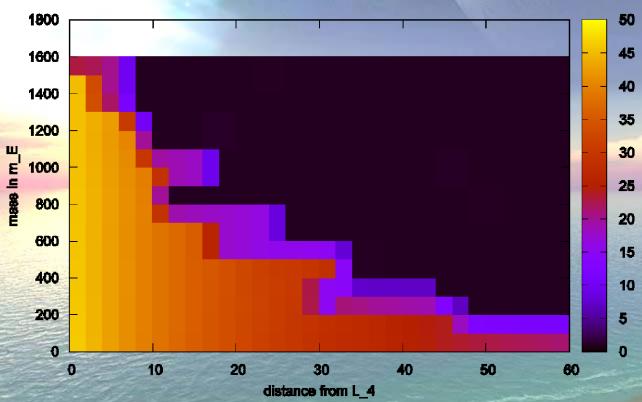


Three equal masses

Extension to very large masses

1500 = 5 M_Jup

JJJ

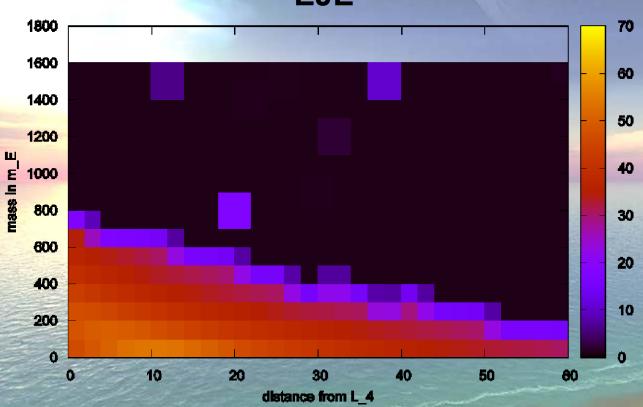


Two equal masses

Extension to very large masses

1500 = 5 M_Jup





Conclusions

Stable region around equilibrium points

$$a=a+/-0.02$$

m<1500 m_E

Libration for L_4 for I<~100deg

Ongoing study:

Mass ratios different

a=a+/-0.0?

Inclinations?